

Free Wave in Vacuum

Dissipative waves in a charge free media

$$\sigma \neq 0 \quad \rho_s = 0$$

$$\vec{J}_e = \sigma \vec{E}$$

$$\vec{\nabla} \times \vec{H} = \sigma \vec{E} + \epsilon \frac{\partial \vec{E}}{\partial t}$$

$$\vec{\nabla} \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}$$

$$\vec{\nabla} \times (\vec{\nabla} \times \vec{E}) = -\frac{\partial}{\partial t} (\vec{\nabla} \times \vec{B})$$

$$= -\mu \frac{\partial}{\partial t} (\vec{\nabla} \times \vec{H})$$

$$= -\mu \sigma \frac{\partial \vec{E}}{\partial t} - \mu \epsilon \frac{\partial^2 \vec{E}}{\partial t^2}$$

$$= \vec{\nabla} (\vec{\nabla} \cdot \vec{E}) - \nabla^2 \vec{E} \quad (\rho_s = 0)$$

$$-\nabla^2 \vec{E} = -\mu \sigma \frac{\partial \vec{E}}{\partial t} - \mu \epsilon \frac{\partial^2 \vec{E}}{\partial t^2}$$

$$= -\mu \epsilon \left(\sigma \frac{\partial \vec{E}}{\partial t} \right) - \mu \epsilon \frac{\partial^2 \vec{E}}{\partial t^2}$$

Let $\vec{E} = \text{Re} [\vec{E}_s e^{-i\omega t}]$

$$E_x(x, y, z, t) = \text{Re} [E_{sx} e^{-i\omega t}]$$

What is \vec{E}_s ?

\vec{E}_s is this vector which when multiplied by $e^{j\omega t}$ and take $\text{Re}[\vec{E}_s e^{-j\omega t}]$ you get \vec{E} ; \vec{E}_s is a function of x, y, z only not of t

$$\begin{aligned}\frac{\partial \vec{E}}{\partial t} &= \frac{\partial}{\partial t} \text{Re}[\vec{E}_s e^{-j\omega t}] \\ &= \text{Re}\left[\frac{\partial}{\partial t} \vec{E}_s e^{-j\omega t}\right] \\ &= \text{Re}[-j\omega \vec{E}_s e^{-j\omega t}]\end{aligned}$$

Maxwell in Phasor form

(going to phasors: $\frac{\partial}{\partial t} \rightarrow -j\omega$)

$$\vec{\nabla} \cdot \vec{D} = \rho$$

$$\vec{\nabla} \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}$$

$$\vec{\nabla} \times \vec{E}_s = -j\omega \mu \vec{H}_s$$

$$\vec{\nabla} \cdot \vec{H} = 0$$

unchanged

$$\vec{\nabla} \times \vec{H} = \vec{J}_c + \frac{\partial \vec{D}}{\partial t}$$

$$\vec{\nabla} \times \vec{H}_s = \vec{J}_{cs} + j\omega \epsilon \vec{E}_s$$

Waves in charge free media in phasor form

$$\rho = 0$$

$$\vec{\nabla} \times \vec{E}_s = -j\omega \mu \vec{H}_s$$

$$\begin{aligned}\vec{\nabla} \times (\vec{\nabla} \times \vec{E}_s) &= -j\omega \mu (\vec{\nabla} \times \vec{H}_s) \\ &= -j\omega \mu (\vec{J}_{cs} + j\omega \epsilon \vec{E}_s)\end{aligned}$$

$$\begin{aligned}
 -\vec{\nabla}^2 \vec{E}_s &= -j\omega\mu(\sigma \vec{E}_s + j\omega\varepsilon \vec{E}_s) \\
 &= (-j\omega\mu\sigma + \omega^2\mu\varepsilon) \vec{E}_s \\
 &= \omega^2\mu(\varepsilon - j\frac{\sigma}{\omega}) \vec{E}_s
 \end{aligned}$$

let $\varepsilon_c = \varepsilon - j\frac{\sigma}{\omega} = \varepsilon \left(1 - j\frac{\sigma}{\omega\varepsilon}\right)$

$$-\vec{\nabla}^2 \vec{E}_s = \omega^2\mu\varepsilon_c \vec{E}_s$$

Note: if $\sigma = 0$ then $\varepsilon_c = \varepsilon$ and $\omega^2\mu\varepsilon$ is real

$$\vec{\nabla}^2 \vec{E}_s + \omega^2\mu\varepsilon_c \vec{E}_s = 0$$

if $\sigma = 0$ then $\varepsilon_c = \varepsilon$
and $\omega^2\mu\varepsilon$ can be solved
for ω^2

Introduce the complex parameter γ

$$\gamma^2 = -\omega^2\mu\varepsilon_c = -\omega^2\mu\left(\varepsilon - j\frac{\sigma}{\omega}\right)$$

$$\vec{\nabla}^2 \vec{E}_s - \gamma^2 \vec{E}_s = 0$$

Solutions are therefore

$$e^{\pm \gamma(\hat{n} \cdot \vec{r})}$$

$$\hat{n} \cdot \vec{r} = n_x x + n_y y + n_z z \quad \text{and} \quad n_x^2 + n_y^2 + n_z^2 = 1$$

$$\frac{\partial^2}{\partial x^2} [E_0 e^{\pm \gamma(\hat{n} \cdot \vec{r})}] + \frac{\partial^2}{\partial y^2} [E_0 e^{\pm \gamma(\hat{n} \cdot \vec{r})}] + \frac{\partial^2}{\partial z^2} [E_0 e^{\pm \gamma(\hat{n} \cdot \vec{r})}] = \gamma^2 E_0 e^{\pm \gamma(\hat{n} \cdot \vec{r})}$$

The type of solution $\sim e^{\pm \gamma(\hat{n} \cdot \vec{r})}$ is called the plane wave because, on a fixed plane determined by $\hat{n} \cdot \vec{r} = \text{constant}$,

\vec{E}_s, \vec{H}_s have constant amplitudes.

For the moment, suppose $\hat{n} \cdot \vec{r} = z$

Some properties:

$$\gamma = \alpha + j\beta \quad \vec{E}_0: \text{constant}$$

The physical field is

$$\begin{aligned} \vec{E}(z, t) &= \text{Re} [\vec{E}_s e^{+j\omega t}] \\ &= \text{Re} [\vec{E}_0 e^{\alpha z} e^{j(\beta z + \omega t)}] \end{aligned}$$

For simplicity, Assume \vec{E}_0 is a real vector

$$= \vec{E}_0 e^{\alpha z} \cos(\beta z + \omega t)$$

Beware: $\vec{E}_s = \text{Re} [\vec{E}_s e^{+j\omega t}]$; this change the wave, it now propagates to the left instead of the right, This in essence changes to polarity of α .

$$t_0 = 0 \quad z = -5$$

$$t_1 = 1 \quad z = -7$$

$e^{\alpha z}$ becomes smaller as z becomes smaller

$$e^{-\alpha z}$$

$$\begin{aligned}\vec{E}(z, t) &= \text{Re}[\vec{E}_0 e^{-\alpha z} e^{j(-\beta z + \omega t)}] \\ &= \vec{E}_0 e^{-\alpha z} \cos(\omega t - \beta z)\end{aligned}$$